Contents lists available at ScienceDirect





Flow Measurement and Instrumentation

journal homepage: http://www.elsevier.com/locate/flowmeasinst

Robust estimators for free surface turbulence characterization: a stepped spillway application



Daniel Valero^{a, c, *}, Hubert Chanson^b, Daniel B. Bung^c

^a IHE Delft Institute for Water Education, Water Resources and Ecosystems, 2611, AX Delft, the Netherlands

^b The University of Queensland, School of Civil Engineering, Brisbane, Australia

^c Aachen University of Applied Sciences (FH Aachen), Hydraulic Engineering Section (HES), Aachen, Germany

ARTICLE INFO

Keywords: Acoustic displacement meter Air-water interface Free surface Statistical theories and models Ultrasonic sensor

ABSTRACT

Robust estimators are parameters insensitive to the presence of outliers. However, they presume the shape of the variables' probability density function. This study exemplifies the sensitivity of turbulent quantities to the use of classic and robust estimators and the presence of outliers in turbulent flow depth time series. A wide range of turbulence quantities was analysed based upon a stepped spillway case study, using flow depths sampled with Acoustic Displacement Meters as the flow variable of interest. The studied parameters include: the expected free surface level, the expected fluctuation intensity, the depth skewness, the autocorrelation timescales, the vertical velocity fluctuation intensity, the perturbations celerity and the one-dimensional free surface turbulence spectrum. Three levels of filtering were utilised prior to applying classic and robust estimators, showing that comparable robustness can be obtained either using classic estimators together with an intermediate filtering technique or using robust estimators instead, without any filtering technique.

Author contribution

Daniel Valero: Conceptualization, Methodology, Software, Validation, Formal analysis, Investigation, Data Curation, Writing - Original Draft, Writing - Review & Editing, Visualization, Hubert Chanson: Conceptualization, Resources, Writing - Review & Editing, Supervision, Project administration, Funding acquisition, Daniel B. Bung: Conceptualization, Writing - Review & Editing, Supervision, Project administration, Funding acquisition.

1. Introduction

Numerous studies have been conducted in the past decades aiming to determine different turbulence properties of the free surface. The main focus has been put on linking free surface dynamics to the hydrodynamics underneath [9,14,17] whereas other studies have tried to address the changes that the free surface impose on the surrounding turbulence [12,16,29]. In the extreme event of air-water flows, turbulence has a dominant role on the free surface break-up [3] and, in the specific case of supercritical flows, the free surface deforms throughout the streamwise direction [11,19,22], and transverse direction [6], reaching unstable configurations, which have been suggested to lead to

self-aeration [32,34].

For spillway flows, different instrumentations have been used to obtain characteristic water levels, including point gauges [21], imaging techniques [4,38]; [41]), acoustic displacement meters [4,11,33,38], phase detection probes [22,30,31,38] and, more recently, Light Detection and Ranging (LiDAR) systems [19]. Other experimental studies have used wire gauges to depict turbulent free surface dynamics, see for instance the hydraulic jump analysis of Mouaze et al. [23] and Murzyn et al. [25]. Other examples of non-intrusive techniques for turbulent free surface characterisation include works of Murzyn and Chanson [24] and Chachereau and Chanson [6]. However, little attention has been put to data filtering although the recordings can include numerous outliers due to faulty signals resulting from challenging flow conditions. Most studies have used the raw data, sampled data without voltage range ends, or data within the bounds defined at +/-3 standard deviations from the mean. Commonly studied flow quantities include mean/median and standard deviation, as well as correlation timescales.

With increasing interest on the water surface dynamics, the focus should be placed on its accurate experimental determination. Different methods can be used to describe the free surface, which can be grouped into:

1. Use of classic estimators (e.g., mean, standard deviation, and others).

* Corresponding author. IHE Delft Institute for Water Education, Water Resources and Ecosystems, 2611, AX Delft, the Netherlands. *E-mail addresses:* d.valero@un-ihe.org (D. Valero), h.chanson@uq.edu.au (H. Chanson), bung@fh-aachen.de (D.B. Bung).

https://doi.org/10.1016/j.flowmeasinst.2020.101809

Received 31 October 2019; Received in revised form 31 July 2020; Accepted 9 August 2020 Available online 5 October 2020 0955-5986/© 2020 Elsevier Ltd. All rights reserved.

Notation			correlation (–)
		$\tilde{R_{max}}$	Maximum cross-correlation, obtained using Spearman's
ADM	Acoustic Displacement Meter		rank correlation (–)
DVED	Depth-Velocity Elliptical Despiking	$R_{\eta\eta}$	Autocorrelation function for η (–)
d	Flow depth time series (m)	$\overline{R_{\eta\eta}}$	Classic autocorrelation function for η (–)
d	Mean flow depth (m)	$\tilde{R_{nn}}$	Autocorrelation function for η by means of Spearman's
Ĩ	Median flow depth (m)		rank correlation (–)
ď	Flow depth expected deviation (m)	ROC	Robust Outlier Cutoff
ď	Flow depth standard deviation (m)	STD	Standard Deviation
Ĩ	Flow depth robust standard deviation ($k \operatorname{MAD}(d)$) (m)	$T_{\eta\eta}$	Turbulent timescale for the free surface fluctuation (s)
$\tilde{d_{r}}^{'}$	Flow depth interguartile range (m)	$\overline{T_{\eta\eta}}$	Turbulent timescale for the free surface fluctuation
$d_{n}^{j'}$	Flow depth relative to percentile n^{th} (m)		obtained using the Pearson's product moment correlation
$\frac{d''}{d''}$	Flow depth expected skewness $(-)$	Ť	Turbulent timescale for the free surface fluctuation
<i>d</i> ″	Flow depth quartile coefficient of skewness $(-)$	1 ηη	obtained using the Spearman's rank correlation (s)
E_{m}	One-dimensional flow depth spectra (m^2s^{-1})	t	Time (s)
	One-dimensional flow depth spectra, classic determination	и	Ranked vector (–)
111	(m^2s^{-1})	v_s	Vertical velocity of the free surface (ms^{-1})
$\tilde{E_{nn}}$	One-dimensional flow depth spectra, robust determination	$\overline{v_{s}}'$	Vertical velocity standard deviation (ms ⁻¹)
-1-1	(m^2s^{-1})	$\overline{v_{s}'}$	Vertical velocity standard deviation (ms^{-1})
f	Frequency (s ⁻¹)	$\tilde{v_{s'}}$	Vertical velocity robust standard deviation ($k \text{ MAD}(v_s)$)
k	Relation between the Standard Deviation and the Median		(ms ⁻¹)
	of the Absolute Deviation	w	Ranked vector (–)
LUV	Lower and Upper Voltage	x	Auxiliary vector (–)
MAD	Median Absolute Deviation	z	Auxiliary vector (–)
MED	Median	η	Time series for the deviation from the median flow depth
N	Number of data points (–)		(m).
PMF	Probability Mass Function	λ_u	Universal threshold (–)
<i>R_{max}</i>	Maximum cross-correlation, obtained using classic		

2. Use of robust estimators (e.g., median, median absolute deviation, and others).

Traditionally, classic estimators have been the default choice. These statistical methods are best for opportune situations where raw data perfectly match physical characteristics. Nonetheless, true data are commonly contaminated by errors inherent to any experimental methodology, and filtering should be prescribed to reduce the impact on the turbulence estimations. Alternatively, robust estimators can be used; understanding "robust" or resistant as the characteristic to be affected only to a limited extent by a number of gross errors [15]. This is achieved when a small subset of the sample cannot have a disproportionate effect on the estimators, although they are best for a broad range of situations, tolerate a large quantity of outliers mixed within the data sample (up the breakdown point) without resulting in a meaningless estimate, and perform superiorly even for small data samples [15].

A simple example can be presented through the estimation of the expected value of a variable. As a dataset contains outliers, the use of the mean (classic) estimator should be preceded by a filtering step. Otherwise, a median (robust) estimator can be used, being more insensitive to the presence of outliers. When outliers are *obvious*, they can still be removed, but accurate determination of the expected value does not strongly rely on the adequacy of the filtering technique.

This work explores the aforementioned dual data analysis by studying different combinations of filtering techniques and classic/ robust estimators for an extremely turbulent flow case: the turbulent free surface in the non-aerated region over a stepped spillway (introduced in Section 2). Robust estimators, mainly based on the simple concept of median and data ranking, are proposed in Section 3 for a wide range of turbulence properties, namely: the expected free surface level, the expected fluctuation intensity, the depth skewness, the autocorrelation timescales, the vertical velocity fluctuation intensity, the perturbations celerity and the one-dimensional free surface turbulence spectrum. Alternatively, three filtering techniques based upon well-stablished works [13,35] are presented in Section 4. These three techniques present gradually increasing intricacy, naturally encompassed by higher rejection rates. Section 5 analyses the combination of classic and robust estimators with the proposed filtering techniques for different turbulence quantities of the turbulent free surface over a stepped spillway. Results discussion and final conclusions are presented in Sections 6 and 7, respectively.

2. Experimental setup

2.1. Geometry

The present study focuses on the non-aerated region of a stepped spillway model of 45° slope (1V:1H) located at The University of Queensland. The stepped spillway has a wide inlet basin (5 m wide, 2 m long) which ensures smooth inlet conditions, leading to a broad crested weir (0.60 m long, 0.985 m wide) which conveys the flow into the stepped spillway (same width) composed of steps of height h = 0.10 m. The water discharge is estimated using a previously obtained experimental discharge relationship based on detailed velocity measurements. A thorough description of this spillway geometry, and other flow variables, can be found in Zhang and Chanson [37]. The experimental setup is shown in Fig. 1.

2.2. Instrumentation

Instantaneous free surface measurements were sampled with three microsonic[™] Acoustic Displacement Meter (ADM) mic+25/IU/TC. The measuring range recommended by the manufacturer is 30–250 mm. The



Fig. 1. Stepped spillway model. (a) Image taken at 1/1000 shutter speed, processed with contrast-limited adaptive histogram equalization to enhance the characteristic free surface perturbations and air entrainment after step IV. (b) Sketch of the experimental setup (rotated 45° counterclockwise), steps numbering, sensors measurements location (- -), pseudobottom and parallel axis (). Flow from left to right.

near field of the ADM sensors was enclosed with PVC cylinders of the same diameter to prevent the wetting of the sensors. This artefact did not alter the sensors' output signal, as shown by Kramer and Chanson [18].

The ADM sensors provide a voltage time series that can be correlated to a distance in order to estimate a water level. The three ADM sensors were calibrated over a distance range covering the expected water depths to be measured. Calibration was conducted by recording during 300 s at 100 Hz for 11 different distance levels, which covered the range of expected flow depths. Fig. 2a shows that calibration exhibited a linear relation over the entire sampled range (note that different placement of the sensors yield different voltage-depth relationships) and Fig. 2b that the Standard Deviation (STD) of each calibration step remained close to 0.10 mm, computed from the voltage STD and using the obtained calibration curve, which matches the accuracy specified by the sensors manufacturer.

2.3. Measurement location and flow conditions

The ADM 1 was located at a fixed position over the crest, 0.17 m upstream over the downstream edge of the broad crest (step 0, Fig. 1b). The other two sensors (ADM 2 and ADM 3) were located over the stepped geometry separated by 0.141 m in the longitudinal direction, thus coinciding with one cavity length. Keeping a constant distance between both sensors, ADM 2 and ADM 3 were placed above the pseudobottom (formed by the step edges, Fig. 1b), allowing the measurement of the flow depths at different spillway locations. Recordings were conducted at the step edges (steps 0 - VII) and above the step cavities (mid distance between the step edges), as marked in Fig. 1. Each

recording was conducted at a sampling rate of 100 Hz during 600 s. The total time recorded and the distribution of the measurements over the spillways is shown in Fig. 3. Differences of sampling time are due to two reasons: overlapping of measurement locations as the ADM sensors were moved downstream; and repetition of some measurements at locations where the free surface was highly roughened due to turbulence. The inception point of air entrainment is marked for each investigated discharge ($d_c/h = 0.9, 1.1, 1.3, 1.5, 1.7, 1.9$ and 2.1, with d_c the critical depth) according to the visual observations of Zhang and Chanson [37]. For reference, the empirical formulas of Meireles et al. [22] and Chanson et al. [8] are included. Precisely, all the measurements fall within the non-aerated region, where the flow gradually roughens as the flow becomes more turbulent up to break up [34].



Fig. 3. Measurements distribution and sampled time for the studied flow conditions. Visual observation of the inception point location [37] and empirical formulas of Meireles et al. [22] and Chanson et al. [8].



Fig. 2. ADM performances. (a) ADM calibration curves (mean flow depth \overline{d} and corresponding mean voltage level) and (b) Depth fluctuation $\overline{d'}$ for static measurements (noise level) corresponding to each mean flow depth level measured during the ADM calibration curve determination. 'Resolution' specified by sensors' manufacturer.

3. Robust estimators

The use of estimators insensitive to the presence of outliers can yield more reliable turbulence predictions, hence alleviating the responsibility often relying solely on the filtering techniques. In the following, the estimators corresponding to classic statistic techniques are presented overlined (e.g., \overline{d}) whereas the robust counterpart is presented with a tilde (e.g., \overline{d}).

3.1. Expected value and fluctuation

The expected value of a variable ($E[\cdot]$) is the first variable of interest in any data analysis. For the case of the flow depths, it is often estimated by using the mean (\overline{d}) , defined as the ensemble average of all the samples in the filtered signal:

$$\mathbf{E}[d] := \overline{d} = \frac{1}{N} \sum_{i=1}^{N} d_i \tag{1}$$

with N the total number of flow depth measurements. Deviations from the expected value of a variable can also be of interest, as they are associated to turbulence and dispersion, and can be studied on the basis of:

$$\eta = d - \mathbf{E}[d] \tag{2}$$

By definition, $E[\eta] = 0$ and it is more frequent to study the expected value of the squared fluctuation instead. An estimation of the dispersion of the data can be done by means of the sample standard deviation (STD) as:

$$\sqrt{\mathrm{E}[\eta^2]} := \overline{d} = \sqrt{\frac{\sum_{i=1}^{N} \left(d_i - \overline{d}\right)^2}{N - 1}}$$
(3)

The STD approximates the population variance using the squared value of each sample deviation, thus endorsing bigger weight to the outliers which depart significantly from the expected value of the series. The mean value can be affected by the presence of outliers as well but, when equally distributed around the mean, their contribution to Eq. (1) would balance. Alternatively, the median (MED) and the Median Absolute Deviation (MAD) can be used as estimators of location and variance, being both robust estimators against outliers with a breakdown point of 50% (i.e., 50% of contaminated data is necessary to force the estimator to result in a false output) as opposed to the counterpart mean and standard deviation, which hold a 0% breakdown point. It is noteworthy that the median is the location estimator that presents the highest breakdown point [20] and is defined as the value separating the greater and lesser halves of the series. Hence, a robust estimator for the expected value is herein proposed directly through the median operator $\tilde{d} = \text{MED}(d)$. Similarly, the MAD represents the best robust scale estimator, even more than the interquartile range that remains at a 25% breakdown point [20,27].

The MAD can be obtained by sorting the absolute value of the residuals around the MED and selecting the value corresponding to the 50%. Nonetheless, it is implemented in many commonly used numerical libraries (e.g., MATLAB®, R programming language or Python 2.7 together with the *statsmodels* library, being the latter combination the one used in this study). The MAD of the sampled flow depth can be related to the standard deviation of different probability density functions as [27]:

$$d' = k \operatorname{MAD}(d) = k \operatorname{MED}(|\eta|)$$
(4)

being k a coefficient related to the sample distribution. When a Gaussian behaviour is assumed, k takes the value [27]:

$$k = 1.483$$
 (5)

Another estimation of the flow depth variance can be obtained through the interquartile range $\tilde{d_1} = \tilde{d_{75}} - \tilde{d_{25}}$, being $\tilde{d_{75}}$ and $\tilde{d_{25}}$ the depth levels representing the 75th and 25th percentiles, respectively. This estimator presents a lower breakdown point than the MAD, therefore being more sensitive to the presence of outliers.

3.2. Skewness

Higher order statistics can be computed to study the shape of the free surface waves. With increasing order, the exponent weighting the outliers is also incremented although for even order numbers some positively and negatively deviated outliers could balance.

The flow depth skewness $(\overline{d''})$ can be defined as.

$$\overline{d''} = \frac{1}{N} \sum_{i=1}^{N} \left(\frac{d_i - \overline{d}}{\overline{d}} \right)^3 \tag{6}$$

This descriptive statistic is a dimensionless measure of the lack of symmetry. Following the robust estimators defined for first and second order statistics, quartiles information can be used to define a robust estimator for the skewness [39]:

$$\tilde{d}'' = \frac{\tilde{d}_{75} - 2\,\tilde{d} + \,\tilde{d}_{25}}{\tilde{d}_{I}'} \tag{7}$$

The parameter \tilde{d}'' is the so-called quartile coefficient of skewness [39].

3.3. Autocorrelation timescales

The cross-correlation function between two series (x and z) can be computed as:

$$R_{xz}(\tau) = \frac{\operatorname{cov}(x, z(\tau))}{\sqrt{\operatorname{E}\left[(x - \operatorname{E}[x])^2\right]}\sqrt{\operatorname{E}\left[(z(\tau) - \operatorname{E}[z])^2\right]}}$$
(8)

being cov the covariance and τ the lag time. Efficient computation of the covariance for long samples can be achieved through fast Fourier transformation. Note that some terms of Eq. (8) can be rewritten using classic estimators as:

$$\overline{R_{xz}}(\tau) = \frac{\operatorname{cov}(x, z(\tau))}{\overline{x'} \overline{z'(\tau)}}$$
(9)

An alternative, nonparametric form of the correlation can be defined by means of the Spearman's correlation. For this purpose, both x and zare ranked separately from smallest to largest values, assigning the mean rank when equal values occur. Let u_i and w_i take the rank of the ith observation in x and z, respectively. Spearman's rank correlation of xand z can then be computed as [39]:

$$\tilde{R}_{xz}(\tau) = \overline{R_{uw}}(\tau) = \frac{\operatorname{cov}(u, w(\tau))}{\overline{u'} \,\overline{w'(\tau)}}$$
(10)

The use of the ranked data has the advantage that it allows computation of the correlation between both trends without strongly depending upon the current value of each measurement. This alternative correlation does not considerably slow down the computation as the ranking of the vectors can be done with efficient sorting algorithms and, afterwards, the sorted data can be correlated as per the original vectors. Spearman's correlation is the nonparametric version of the Pearson correlation coefficient, which makes its computation more robust to outliers.

Cross-correlation function can be used to obtain the most probable lag between two time series $[= \operatorname{argmax}(R_{xz})]$ or to obtain the autocorrelation function (R_{xx}) , that allows extraction of turbulent scales and spectrum. A timescale $(T_{\eta\eta})$ for the flow depth fluctuations can be obtained by integrating its autocorrelation function ($R_{\eta\eta}$) up to the first zero-crossing point:

$$T_{\eta\eta} = \int_{0}^{\tau=\tau(R_{\eta\eta}=0)} R_{\eta\eta}(\tau) \,\mathrm{d}\tau \tag{11}$$

Depending on how the autocorrelation function is computed (Eq. (9) or 10), a classic estimation of the turbulent timescale $(\overline{T_{\eta\eta}})$ could be carried out through $\overline{R_{\eta\eta}}$ or, an alternative, robust turbulent timescale $(\tilde{T_{\eta\eta}})$ could be computed through $\tilde{R_{\eta\eta}}$.

3.4. Cross-correlation peak and wave celerity

Given that two ADMs are placed in series (ADM 2 and ADM 3), their signals can be cross-correlated allowing estimation of the most probable time lag. When the distance between the synchronized sensors is known, the celerity of the free surface perturbations can be estimated. Both the classic correlation (Eq. 9) and the Spearman's correlation (Eq. 10) can be used to compute the cross-correlation and, hence, estimate the waves' celerity (\bar{c} and \tilde{c} , respectively).

3.5. One-dimensional flow depth spectrum

The one-dimensional spectrum has been traditionally computed for velocity time series. It allows insight into the flow structure and the energy distribution for different wavelengths (i.e., different eddy sizes). In this study, the one-dimensional spectrum, as defined by Pope [26]; is proposed for the flow depth fluctuation:

$$E_{\eta\eta}(f) = \frac{2}{\pi} E[\eta^2] \int_0^\infty R_{\eta\eta}(s) \cos(f s) \, \mathrm{d}s$$
 (12)

Both $E[\eta^2]$ and $R_{\eta\eta}$ can be estimated either in a classic (Eqs. 3 and 9) or a robust manner (Eqs. 4 and 10), thus leading to a standard ($\overline{E_{\eta\eta}}$) or robust ($\tilde{E_{\eta\eta}}$) estimation of the one-dimensional depth fluctuation spectrum.

4. Data filtering

A data point is oftentimes labelled as an outlier when it lies at an abnormal distance from other values of a certain population. Nonetheless, for any observation far from the group, there is a positive (despite small) possibility to occur and thus the crux is on identifying these outliers without losing true information from the population [1]. Doubtful or anomalous values can come from a mixed sample of a different population or erroneous measurements. It is also important to understand how outliers are physically generated to understand their likelihood of occurrence.

4.1. Lower and upper voltage (LUV)

When an ADM pulse-echo is lost, the sensor is incapable of generating a proper estimation of the free surface position. The voltage provided for these lost echoes usually piles at the highest or lowest voltages, far away from the voltage values corresponding to realistic depths. It is therefore convenient to locate the sensor so that the measurements are contained in a region of interest far from the extreme voltage values. A first filtering approach could be to simply remove values below and above 5% and 95% from the total voltage range of the ADM. This double threshold filtering technique is based on a physical observation. For the ADM model used in this study, the voltage filtering levels correspond to 0.5 V and 9.50 V respectively (Fig. 4a).

4.2. Robust outlier cutoff (ROC)

Outliers depart from the expected estimation of the flow depth, but do not necessarily accumulate out of the LUV bounds. The Probability Mass Function (PMF) shows that some erroneous measurements run together at different voltage levels. A quick flow observation indicates that these voltage values, associated to different water levels, are not physically meaningful and the introduction of narrower bounds arises as a preferred alternative than LUV filtering.

A commonly used technique is to estimate the variance of the sample, by means of the STD, to establish the filtering bounds around a certain number of STD away from the mean. An alternative way to estimate variance can be done through robust estimators, as presented in Eq. (4). Difference between normal and robust estimators can be wellperceived in Fig. 4b, where a Gaussian function is fitted using the location and variance obtained with the classic estimators (mean and STD) and through the robust estimators (MED and MAD). Fig. 4b also shows a small proportion of outliers piling up at different voltage levels, which are readily observable when using a vertical log-scale.

On the question of how many standard deviations are necessary to be accounted for to make sure that "good data" is not filtered out, the universal threshold represents a conservative estimator. It can be expressed as [13]:

$$\lambda_u = \sqrt{2\ln(N)} \tag{13}$$

with *N* the total number of data points of the sample. Use of the universal threshold yields bounds wide enough to avoid filtering out good data, even if the underlying distribution is slightly skewed, but (usually) narrower bounds than those proposed by the LUV technique. If the final



Fig. 4. First 600 s of voltage signal of ADM 3, $d_c/h = 2.1$ and step V. The number of outliers removed using LUV technique corresponds to 1.3% of the total number of samples. (a) Sensor's signal (temporal series); (b) voltage levels distribution and Gaussian fit using classic (normal) and robust estimators.

distribution is markedly skewed, MAD can be estimated for both positive and negative deviations departing from the MED value and, consequently, different filtering thresholds could be defined for positive and negative deviations.

4.3. Depth-velocity elliptical despiking (DVED)

One step further on the filtering of the flow depth time series could be conducted using the finite differences of η and its variance, following Goring and Nikora [13] method for velocity data. Provided that an ADM can measure the time series of flow depth, a vertical velocity (v_s) can be estimated by using the central finite difference:

$$v_s \equiv \frac{\partial}{\partial t} = \frac{\partial}{\partial t} \approx \frac{\Delta \eta_j}{\Delta t_j} = \frac{\eta_{j+1} - \eta_{j-1}}{t_{j+1} - t_{j-1}}$$
(14)

For the data at the beginning and ending of the sample, backward or forward differences can be taken by simply using *j* instead of j+1 or j-1 at the right-hand side term of Eq. (14).

For η , maximum and minimum thresholds can be established based on Eqs. (3) and (14) for v_s as well. This approach is a simplified version of the work of Goring and Nikora [13] and Wahl [35]; which extended the analysis up to the second derivative of the variable under analysis. It must be noted that the filtered data could still produce unrealistic vertical velocities, because some depth outliers could randomly fall inside the bounds defined by the ROC technique, hence remaining undetected.

Goring and Nikora [13] proposed ellipsoid-type bounds based on the observation that "good data" tend to cluster together forming this shape. In this study, the 2D PMF for η and v_s of the data was analysed and similar clustering forms were recognized. For conciseness, only the 2D PMF of the data previously shown in Fig. 4a and 4b is presented in Fig. 5. The isoprobability contours (points with same probability of occurrence) appear to take the form of ellipse-like curves. Accounting for the universal threshold as a situation of equal likelihood of appearance, the hypothesis of Goring and Nikora [13] remains consistent for flow depths measurements. Thus, the filtering bounds can be finally expressed as:

$$\left(\frac{\eta}{MAD(\eta)}\right)^2 + \left(\frac{v_s}{MAD(v_s)}\right)^2 \le \left(\lambda_u \ k \ \right)^2 \tag{15}$$

Equation (15) reduces to the ROC filtering when the second term of the left-hand side is neglected.



Fig. 5. Probability Mass Function (PMF) of the data shown in Fig. 4.

4.4. Practical implementation considerations

The three proposed filtering methods correspond to: one physically based filtering (LUV), and two statistical techniques (ROC and DVED). The ROC and DVED algorithms produce upper and lower bounds based on the percentiles of the sampled data, while the LUV method defines bounds based on the physical observation that extremely high and low voltages correspond to lost echoes. Thus, LUV can be complementary to ROC and DVED techniques under certain conditions. An example can be given by a recording for which the free surface is considerably sloped relative to the axis of the ADM resulting on more than 50% of erroneous data. In such case, most measurements would fall close to 10 V and the application of ROC or DVED alone would not result in adequate identification of the outliers. Hence, it is proposed that, after applying ROC and DVED, the LUV technique should be used to avoid accepting erroneous data. In the following, any result presented for ROC and DVED techniques has been also filtered using LUV afterwards. Additionally, if a filtering technique flagged more than 50% as invalid data, the recording was dismissed and not accounted for in the subsequent analysis.

When an outlier is detected by any of the presented filtering methods, it can be simply deleted or replaced. The most basic method would consider the removal of the outlier. The outlier could also be substituted by the average or the median of the entire signal, which would imply that the value replacing the outlier simply takes the "expected" value of the signal. Nonetheless, this can create fast gradients which were not contained in the original signal and, it could considerably affect the latter computed statistical estimators (e.g., by reducing the STD). Another alternative would be using the neighbouring value (sample-and-hold), which presents certain advantages when it comes to spectral representation [28]; [40]). In this study, it is proposed to substitute the outliers by linear interpolation between their surrounding points. It must be noted that linear interpolation replacement strategy can affect slopes of a power spectrum. Higher order polynomials could lead to new spikes [5]. More complex methodologies could be proposed to generate the outlier replacement. Herein, linear interpolation is adopted given that it is the most commonly used approach.

5. Results

5.1. Rejection rate

The three proposed filtering techniques were applied to the data presented in Fig. 3, obtaining different rates of rejection for different measuring locations (see Fig. 6 and Table 1). Fig. 6 shows that large percentages of data were rejected close to the spillway crest (step 0), where free surface bends following the chute axis at the transition from the broad crested weir to the spillway (Fig. 1). This can be explained by the inclination of the detection zone axis of the ADM with respect to the normal to the free surface [38]. These large rejection rates close to the first step were obtained for all discharges, indistinctly of the filtering technique.

With increasing discharge, the flow depth becomes more parallel to the pseudobottom, as opposed to the flow depth curvatures that encompass the step edges that are observed for the lower discharges. Hence, the free surface tends to be closer to the axis of the ADM measuring cone at large discharges and fewer outliers can be expected. Close to the inception point of air entrainment, the free surface considerably roughens with dynamic interfacial processes [7,34] and its dynamic determination can be more challenging for the ADM sensors, consequently resulting in a local increase of the outliers contained in the recorded dataset.

Generally, no major change in the amount of rejected data occurred when applying the ROC technique compared to the LUV method, but approximately twice more data was removed when applying the DVED technique (see Table 1).



Fig. 6. Percentage of rejected data from ADM 2 after application of different filtering techniques: (a) LUV, (b) ROC and (c) DVED.

Table 1 Median percentage of rejected data through the non-aerated region of the spillway.

	Filtering method			
ADM	LUV (%)	ROC (%)	DVED (%)	
2	3.1	3.2	8.3	
3	5.7	6.2	12.2	

5.2. Expected value of the flow depth and its fluctuation

The first analysed variable corresponds to the expected value of the flow depth. A similar prediction with the mean estimator (Eq. 1) is obtained independently of the filtering method used. Positive and negative deviations compensate, resulting in a negligible effect on the mean estimation for almost all the data. For the median estimator, all data points yield an exact same estimation, independently of the filtering technique used. Consequently, both mean and median estimators can provide with accurate values of the expected flow depth regardless of the filtering technique.

In terms of the mean fluctuation of the flow depth (Eq. 3), a

significant deviation from the prediction of \overline{d} is obtained by using the filtered data series using the LUV method as well as the method using the filtered data series of the ROC and the DVED methods (Fig. 7a). In Fig. 7a, both the ROC and the DVED filtered data converged to similar values, with differences below 10%. It is then proposed that when studying the depth variance, at least the ROC method should be used, instead of the LUV method, to avoid incorporating outliers as an unphysically higher turbulence level. An alternative approach, if the analysis is just restricted to mean and mean fluctuation levels, is to filter the data using the LUV method but taking advantage from the robust behaviour of the MAD estimator (Eq. 4) to approximate the samples' expected fluctuation. In such case, comparison of \tilde{d} estimation between the LUV and the DVED filtered data (Fig. 7b) revealed that roughly all the data falls between perfect agreement and -10% lines, as it occurs for the ROC and DVED filtered data for $\overline{d'}$ (Fig. 7a). Performance for $d_{I'}$ was in close agreement to that of \tilde{d}' (shown in Fig. 7b), with almost all the predictions varying less than 10% indistinctly of the filtering method.



Fig. 7. Effect of different filtering techniques on different turbulent flow variables using traditional (left) and robust (right) estimators. (a, b) expected value of the depth fluctuation, (c, d) skewness and (e, f) turbulent timescales. Note that for (a–f) ROC and DVED perform likewise and that (c) does not include LUV filtered data (which showed random, out of order spreading). Perfect agreement (-), $\pm 10\%$ deviation (--) and $\pm 20\%$ deviation (-·).

5.3. Flow depth skewness

Discrepancies between $\overline{d''}$ estimated based upon data obtained after filtering with the ROC and the DVED algorithms and the classic skewness estimator (Eq. 6) remained usually below 20%. Using only the LUV filtering method, the remaining outliers scaled up by two orders of magnitude the skewness predictions. This is shown in Fig. 7c, as a comparison between the results obtained after filtering with the ROC and the DVED algorithms and the classic skewness estimator (Eq. 6).

Fig. 7d shows the comparison for all three filtering techniques using robust estimators (Eq. 7). It must be noted that $\tilde{d''}$ is restricted to values between -1 and +1, which prevents from direct comparison to $\overline{d''}$. Nonetheless, a similar trend can be observed for both classic and robust estimators, with the latest allowing reasonable skewness estimations even for the less restrictive filtering technique (LUV).

5.4. Autocorrelation timescales

Autocorrelation timescales estimated with the classic estimators (Eqs. 9 and 11) are considerably smaller when the LUV method is applied, instead of the ROC or DVED (Fig. 7e). Likewise, data filtered with the ROC technique led to predictions 20% smaller for $T_{\eta\eta}$ than

when using the DVED method (Fig. 7e). Results for the robust estimation of the autocorrelation timescale ($T_{\eta\eta}$, Eqs. 10 and 11) are shown in Fig. 7f. Using the robust Spearman's ranked autocorrelation, the differences in the estimation of the autocorrelation timescales reduce significantly between the data filtered using the LUV method and the other more stringent filtering techniques. Fig. 7e and 7f shows that the effect of noise in the sampled data is to reduce the autocorrelation function values, which yields estimations corresponding to shorter or faster eddies. For all cases, the numerical integration of Eq. (11) was conducted using the trapezoidal rule.

5.5. Vertical velocity fluctuation

An estimation of the instantaneous free surface vertical velocity can be obtained using Eq. (14), with a zero-mean value in a steady flow. The intensity of the vertical velocity fluctuation can be studied in terms of the STD of the velocity time series ($\overline{v_s'}$). The resulting $\overline{v_{s'}}$ values obtained from the filtered data of all three proposed techniques follow closely the same trend as results for \overline{d} . The data filtered using the LUV method produces considerably larger values of $\overline{v_{s'}}$ than the other two filtering techniques, as some noise is incorporated as a turbulence level. When taking advantage of the robust nature of the MAD to estimate $v_{s'}$, similar



Fig. 8. Effect of different filtering techniques on the free surface perturbations celerity obtained with: (a) traditional cross-correlation (\bar{c}) and (b) with Spearman's ranked cross-correlation (\tilde{c}). Perfect agreement (–), ±10% deviation (- -) and ± 20% deviation (--).

results for all three filtering techniques are obtained indistinctly of the filtering method; likewise for \tilde{d} , shown in Fig. 7b. Hence, discrepancies in the robust estimation between the LUV and DVED filtered data remain around +20% deviation and the perfect agreement lines.

5.6. Free surface perturbation celerity

Fig. 8a shows that the data filtered using the two most stringent methods, ROC and DVED methods, yields similar \overline{c} estimations. Differently, estimations based on the LUV filtered data scatter significantly. Maximum correlation ($\overline{R_{max}}$) is also shown in Fig. 8a as it is an indicator of similarity between the two cross-correlated signals. The median value for all measurements of $\overline{R_{max}}$ was 0.16 for the LUV filtered data, 0.45 for the ROC filtered data and 0.44 for the DVED filtered data; showing a clear improvement with the most restrictive filtering techniques.



Fig. 9. One-dimensional free surface spectrum based on: (a) standard estimators $(\overline{E_{\eta\eta}})$ or (b) robust estimators $(\overline{E_{\eta\eta}})$; all averaged over 60 non-overlapping spectra. Data corresponding to one recording of 600 s of ADM 2 for step I-II (cavity) and $d_c/h = 1.1$.

Using the Spearman's based cross-correlation method (Eq. 10), the estimations generally relied between the +/-20% accuracy range, even for the LUV and DVED filtered data (Fig. 8b). Both the ROC and the DVED filtered data estimations coincided for most measurements. The maximum correlation R_{max} was 0.42 for the LUV filtered data, 0.44 for the ROC filtered data and 0.41 for the DVED filtered data; showing similar levels than after applying the ROC filtering technique with the classic correlation approach.

5.7. One-dimensional flow depth fluctuation spectrum

Spectrum shown in Fig. 9 were obtained by dividing the 600 s samples into 60 equal length non-overlapping signals; which is long enough based on results of Fig. 7e and 7f for the autocorrelation time-scales that held values generally around 0.5 s. The resulting spectra were subsequently ensemble-averaged. This procedure is based on that proposed by Welch [36]; despite the temporal window has been chosen arbitrarily large as to comprehend a wide range of turbulent timescales. The one-dimensional spectra were obtained for different flow conditions (Fig. 3) but, for the sake of briefness, the effect of the three filtering techniques and the use of classic and robust estimators is only shown for one location and flow condition, albeit similar conclusions can be derived from the others.

Fig. 9a shows that ROC and DVED filtering techniques together with classic estimators yield similar power law slopes and energy levels,

Table 2

Coefficient of determination [2] for different turbulence quantities. Combination of traditional indicators with raw data, LUV and ROC filtered data. DVED filtered data used as reference.

	Variable estimated					
Filtering	d	ď	$\overline{d''}$	$\overline{T_{\eta\eta}}$	$\overline{\nu_{s}'}$	ī
Raw	0.978	0.745	-0.278	0.458	0.724	0.441
LUV	1.000	0.767	-0.153	0.750	0.746	0.697
ROC	1.000	1.000	0.984	0.996	0.998	0.981

Table 3

Coefficient of determination [2] for different turbulence quantities. Combination of robust indicators with raw data, LUV and ROC filtered data. DVED filtered data used as reference.

	Variable estimated						
Filtering	Ĩ	đ	$\widetilde{d'_I}$	<i>đ</i> ″	$T ilde{\eta}\eta$	$\tilde{\nu_{s'}}$	ĩ
Raw	1.000	0.993	0.986	0.839	0.070	0.972	0.484
LUV	1.000	1.000	1.000	0.995	0.845	0.996	0.655
ROC	1.000	1.000	1.000	0.997	0.996	0.996	0.787

although the application of the LUV technique results in a flatter spectrum at lower energies and considerably larger higher energies at all frequencies – associated to considerably higher error levels – as expected from the results for the flow depth fluctuation estimation. Fig. 9b shows that all three spectra based upon robust estimators are in very close agreement while the power slopes are maintained and coincide with those theoretically obtained by Valero and Bung [34]. This holds true for early stages of perturbations development, for large displacements the gravitation waves (low frequencies) flatten (see Ref. [32] for further insights on flow process). This result highlights that robust techniques allow accurate determination of the turbulence spectrum independently of the employed filtering method.

6. Discussion

The application of robust estimators, as opposed of classic estimators, yield results not directly based on the value of each flow quantity measured, but on its distribution (i.e., percentiles information). When using a robust estimator for the study of a flow variable, the exact value of an outlier does not alter the turbulence estimation but only moves a percentile position, which is later retrieved (or not, depending on the position of the outlier) for the turbulence estimation. Consequently, the change in the final turbulence estimation is considerably smaller.

Random errors may be also included within the sampled data, in the form of electrical noise for example, and these might also affect turbulence estimations. When using classic estimators, the statistical moment used to describe a turbulence quantity imposes the same order over the noise deviations. For the robust estimator counterpart, the random noise imposes a certain diffusion in the PMF that is transferred to the turbulence estimations but, nonetheless, does not scale with the power of the statistical moment.

The main limitation of robust estimators comes from the assumption of an underlying probability distribution for a certain variable. Turbulence may well not be Gaussian [10], and differences should be incorporated for instance in Eq. (5) to enable more refined flow's turbulence parametrizations. If departure from Gaussianity is smaller than the errors affecting experimental sampling, robust estimators may by default be more accurate. Besides, filtering of high frequencies may considerably reduce the impact of random noise, but frequency domain filtering fell out of the scope of this investigation.

A summary of the filtering techniques' performance and the effect of robust estimators on most of the studied variables is presented in Tables 2 and 3 and in Fig. 10. The coefficient of determination, defined as





Fig. 10. Coefficient of determination for different flow variables using the raw data $(- \cdot -)$, LUV filtered data $(- \cdot)$ and ROC filtered data $(- \cdot)$ against the DVED filtered data. (a) Traditional estimators and (b) alternative robust estimators.

the squared value of the Pearson's product moment [2], is used to assess the efficiency of the proposed techniques. This coefficient ranges from 0 to 1, with 0 for no correlation and 1 for perfect correlation. Nonetheless, this efficiency estimator cannot detect bias, which is better observed in Figs. 7 and 8. Raw data has been also included in Fig. 10 for completeness.

In Fig. 10, the performance is defined against the DVED filtered data, being the most restrictive filtering technique. Nonetheless, the ROC technique achieves similar results with just half the amount of the rejected data (see Fig. 6). Both the raw and the LUV filtered data can yield similarly accurate estimations of expected depths and variance, when used together with robust estimators (MED and MAD). For the skewness determination, the LUV filtering method and robust estimators are the minimum data processing level which should be used. Nevertheless, a similar degree of complexity is involved when using the ROC method and the amount of filtered data does not increase considerably. For the one-dimensional spectrum, the proposed robust non-parametric method allowed detection of the power law scaling whereas power levels converged even with lowest levels of filtering. Robust estimators

should be used when possible to reduce the uncertainty of the turbulence predictions, as they are insensitive to the presence of outliers. The only exception observed in this study is the case of the waves' celerity determination. It is here hypothesized that small differences in an instantaneous depth value, can produce a large difference of rank within a large signal. These, due to random errors, are different in both signals which can result in two ranks with larger differences, thus impairing their cross-correlation.

7. Conclusions

Classic and robust estimators together with three different filtering techniques have been investigated aiming to shed some light on the accurate determination of free surface turbulent quantities. Measurements were conducted in the non-aerated region of a large stepped spillway model (Fig. 1), which represents one of the most challenging turbulent free surface flows investigated in literature. The performances of the investigated classic and robust estimators and filtering techniques were assessed for the most common turbulence flow variables, ranging from simple fluctuation intensities to the autocorrelation timescales and one-dimensional turbulence spectrum.

The findings show that:

- When using classic estimators, filtering based on the ROC method is a minimum to provide sound turbulence estimations.
- The results did not significantly change with the application of the most restrictive filtering method (DVED), although the amount of rejected data was doubled.
- Robust estimators can provide an alternative to classic estimators while being more insensitive to the presence of outliers.

For all studied turbulent variables, at least a resilient estimator has been proposed. Using robust estimators, the turbulence predictions remained accurate even when the filtering technique would be insufficient with classic estimators. Whereas some estimators hold parallelism or equivalence (see mean/median STD/MAD), others do not necessarily satisfy that, but aim to characterise the same variable feature. Using the robust alternative can help identifying flow features across different studies, using different instrumentations subject to different levels of error.

Funding

The presented results were obtained in the framework of a projectrelated personal exchange program funded by the German Academic Exchange Service (DAAD) and Australia-Germany Joint Research Cooperation scheme with financial support of the Federal Ministry of Education and Research (BMBF).

7.1. Data availability

Ultrasonic sensors recordings and calibration curves are available under: https://doi.org/10.17632/ndx8wkrfb5.1.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

References

- [1] V. Barnett, T. Lewis, Outliers in Statistical Data, John Wiley & Sons, 1978.
- [2] N.D. Bennett, B.F. Croke, G. Guariso, J.H. Guillaume, S.H. Hamilton, A.J. Jakeman, V. Andreassian, Characterising performance of environmental models, Environ. Model. Software 40 (2013) 1–20, https://doi.org/10.1016/j.envsoft.2012.09.011.

- [3] M. Brocchini, D.H. Peregrine, The dynamics of strong turbulence at free surfaces. Part 1. Description, J. Fluid Mech. 449 (2001) 225–254, https://doi.org/10.1017/ S0022112001006012.
- [4] D.B. Bung, Non-intrusive detection of air-water surface roughness in self-aerated chute flows, J. Hydraul. Res. 51 (3) (2013) 322–329, https://doi.org/10.1080/ 00221686.2013.777373.
- [5] L. Cea, J. Puertas, L. Pena, Velocity measurements on highly turbulent free surface flow using, ADV. *Experiments in Fluids* 42 (3) (2007) 333–348, https://doi.org/ 10.1007/s00348-006-0237-3.
- [6] Y. Chachereau, H. Chanson, Free-surface fluctuations and turbulence in hydraulic jumps, Exp. Therm. Fluid Sci. 35 (6) (2011) 896–909, https://doi.org/10.1016/j. expthermflusci.2011.01.009.
- [7] M.R. Chamani, Air Inception in Skimming Flow Regime over Stepped Spillways, in: H.E. Minor, W.H. Hager (Eds.), Air Inception in Skimming Flow Regime over Stepped Spillways. *International Workshop on Hydraulics of Stepped Spillways*, Zürich, Switzerland, Balkema Publ., 2000, pp. 61–67.
- [8] H. Chanson, D.B. Bung, J. Matos, Stepped spillways and cascades, in: H. Chanson (Ed.), Energy Dissipation in Hydraulic Structures vol. 3, CRC Press, Taylor & Francis Group, Leiden, Netherlands, 2015, pp. 45–64.
- [9] D. Dabiri, On the interaction of a vertical shear layer with a free surface, J. Fluid Mech. 480 (2003) 217–232, https://doi.org/10.1017/S0022112002003671.
- [10] M. Farge, K. Schneider, N. Kevlahan, Non-Gaussianity and coherent vortex simulation for two-dimensional turbulence using an adaptive orthogonal wavelet basis, Phys. Fluids 11 (8) (1999) 2187–2201, https://doi.org/10.1063/1.870080.
- [11] S. Felder, H. Chanson, Air-water flows and free-surface profiles on a non-uniform stepped chute, J. Hydraul. Res. 52 (2) (2014) 253–263, https://doi.org/10.1080/ 00221686.2013.841780.
- [12] O. Flores, J.J. Riley, A.R. Horner-Devine, On the dynamics of turbulence near a free surface, J. Fluid Mech. 821 (2017) 248–265, https://doi.org/10.1017/ jfm.2017.209.
- [13] D.G. Goring, V.I. Nikora, Despiking acoustic Doppler velocimeter data, J. Hydraul. Eng. 128 (1) (2002) 117–126, https://doi.org/10.1061/(ASCE)0733-9429(2002) 128:1(117).
- [14] X. Guo, L. Shen, Interaction of a deformable free surface with statistically steady homogeneous turbulence, J. Fluid Mech. 658 (2010) 33–62, https://doi.org/ 10.1017/S0022112010001539.
- [15] D.C. Hoaglin, F. Mosteller, J.W. Tukey, Understanding Robust and Exploratory Data Analysis, John Wiley & Sons, 1983.
- [16] J.C.R. Hunt, J.M.R. Graham, Free-stream turbulence near plane boundaries, J. Fluid Mech. 84 (2) (1978) 209–235, https://doi.org/10.1017/ S0022112078000130.
- [17] E.D. Johnson, E.A. Cowen, Remote monitoring of volumetric discharge employing bathymetry determined from surface turbulence metrics, Water Resour. Res. 52 (3) (2016) 2178–2193, https://doi.org/10.1002/2015WR017736.
- [18] M. Kramer, H. Chanson, Free-surface instabilities in high-velocity air-water flows down stepped chutes, in: Proc. 7th International Symposium on Hydraulic Structures, vol. 15, 2018, p. 18, https://doi.org/10.15142/T3XS8P. Aachen, Germany May 2018.
- [19] M. Kramer, H. Chanson, S. Felder, Can we improve the non-intrusive characterization of high-velocity air-water flows? Application of LIDAR technology to stepped spillways, J. Hydraul. Res. 58 (2) (2020) 350–362, https:// doi.org/10.1080/00221686.2019.1581670.
- [20] C. Leys, C. Ley, O. Klein, P. Bernard, L. Licata, Detecting outliers: do not use standard deviation around the mean, use absolute deviation around the median, J. Exp. Soc. Psychol. 49 (4) (2013) 764–766, https://doi.org/10.1016/j. jesp.2013.03.013.
- [21] I. Meireles, J. Matos, Skimming flow in the nonaerated region of stepped spillways over embankment dams, J. Hydraul. Eng. 135 (8) (2009) 685–689, https://doi. org/10.1061/(ASCE)HY.1943-7900.0000047.
- [22] I. Meireles, F. Renna, J. Matos, F. Bombardelli, Skimming, nonaerated flow on stepped spillways over roller compacted concrete dams, J. Hydraul. Eng. 138 (10) (2012) 870–877, https://doi.org/10.1061/(ASCE)HY.1943-7900.0000591.
- [23] D. Mouazé, F. Murzyn, J.R. Chaplin, Free surface length scale estimation in hydraulic jumps, J. Fluid Eng. 127 (6) (2005) 1191–1193, https://doi.org/ 10.1115/1.2060736.
- [24] F. Murzyn, H. Chanson, Free-surface fluctuations in hydraulic jumps: experimental observations, Exp. Therm. Fluid Sci. 33 (7) (2009) 1055–1064, https://doi.org/ 10.1016/j.expthermflusci.2009.06.003.
- [25] F. Murzyn, D. Mouazé, J.R. Chaplin, Air–water interface dynamic and free surface features in hydraulic jumps, J. Hydraul. Res. 45 (5) (2007) 679–685, https://doi. org/10.1080/00221686.2007.9521804.
- [26] S.B. Pope, Turbulent Flows, Cambridge University Press, 2000.
- [27] P.J. Rousseeuw, C. Croux, Alternatives to the median absolute deviation, J. Am. Stat. Assoc. 88 (424) (1993) 1273–1283, https://doi.org/10.1080/ 01621459.1993.10476408.
- [28] L. Simon, J. Fitzpatrick, An improved sample-and-hold reconstruction procedure for estimation of power spectra from LDA data, Exp. Fluid 37 (2) (2004) 272–280, https://doi.org/10.1007/s00348-004-0814-2.
- [29] M.A.C. Teixeira, S.E. Belcher, Dissipation of shear-free turbulence near boundaries, J. Fluid Mech. 422 (2000) 167–191, https://doi.org/10.1017/ S002211200000149X.
- [30] L. Toombes, H. Chanson, Surface waves and roughness in self-aerated supercritical flow, Environ. Fluid Mech. 7 (3) (2007) 259–270, https://doi.org/10.1007/ s10652-007-9022-y.

D. Valero et al.

- [31] L. Toombes, H. Chanson, Interfacial aeration and bubble count rate distributions in a supercritical flow past a backward-facing step, Int. J. Multiphas. Flow 34 (5) (2008) 427–436, https://doi.org/10.1016/j.ijmultiphaseflow.2008.01.005.
- [32] D. Valero, On the Fluid Mechanics of Self-Aeration in Open Channel Flows, PhD thesis, Université de Liège, Liège, Belgium, 2018, http://hdl.handle.net/2268/22 9191.
- [33] D. Valero, D.B. Bung, Development of the interfacial air layer in the non-aerated region of high-velocity spillway flows. Instabilities growth, entrapped air and influence on the self-aeration onset, Int. J. Multiphas. Flow 84 (2016) 66–74, https://doi.org/10.1016/j.ijmultiphaseflow.2016.04.012.
- [34] D. Valero, D.B. Bung, Reformulating self-aeration in hydraulic structures: turbulent growth of free surface perturbations leading to air entrainment, Int. J. Multiphas. Flow 100 (2018) 127–142, https://doi.org/10.1016/j. ijmultiphaseflow.2017.12.011.
- [35] T.L. Wahl, Discussion of "despiking acoustic Doppler velocimeter data" by Derek G. Goring and Vladimir I. Nikora, J. Hydraul. Eng. 129 (6) (2003) 484–487, https:// doi.org/10.1061/(ASCE)0733-9429(2003)129:6(484).

- [36] P. Welch, The use of fast Fourier transform for the estimation of power spectra: a method based on time averaging over short, modified periodograms, IEEE Trans. Audio Electroacoust. 15 (2) (1967) 70–73.
- [37] G. Zhang, H. Chanson, Hydraulics of the developing flow region of stepped spillways. I: physical modeling and boundary layer development, J. Hydraul. Eng. 142 (7) (2016), 04016015, https://doi.org/10.1061/(ASCE)HY.1943-7900.0001138.
- [38] G. Zhang, D. Valero, D.B. Bung, H. Chanson, On the estimation of free-surface turbulence using ultrasonic sensors, Flow Meas. Instrum. 60 (2018) 171–184, https://doi.org/10.1016/j.flowmeasinst.2018.02.009.
- [39] D. Zwillinger, S. Kokoska, Standard Probability and Statistical Tables and Formulae, Chapman & Hall, CRC Press, 2000, ISBN 1-58488-059-7.
- [40] M. Parsheh, F. Sotiropoulos, F. Porté-Agel, Estimation of power spectra of Acoustic-Doppler velocimetry data contaminated with intermittent spikes, J. Hydraulic Eng. 136 (6) (2010) 368–378, https://doi.org/10.1061/(ASCE)HY.1943-7900.0000202.
- [41] W. Wei, W. Xu, J. Deng, Z. Tian, F. Zhang, Bubble formation and scale dependence in free-surface air entrainment, Sci. Rep. 9 (1) (2019) 11008, https://doi.org/ 10.1038/s41598-019-46883-5.